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Eccentric millisecond pulsars by resonant convection

by

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Red Giant evolution

- Mass-radius relation $r_{\text{env}} \propto \left(\frac{L}{T_{\text{eff}}^4} \right)^{1/2} \propto m_{\text{core}}^{9/2}$
- Mass-period relation $P \propto \left(m_{\text{core}}^{-1/3} r_{\text{env}} \right)^{3/2} \propto m_{\text{core}}^{25/4} \approx m_{\text{core}}^6$

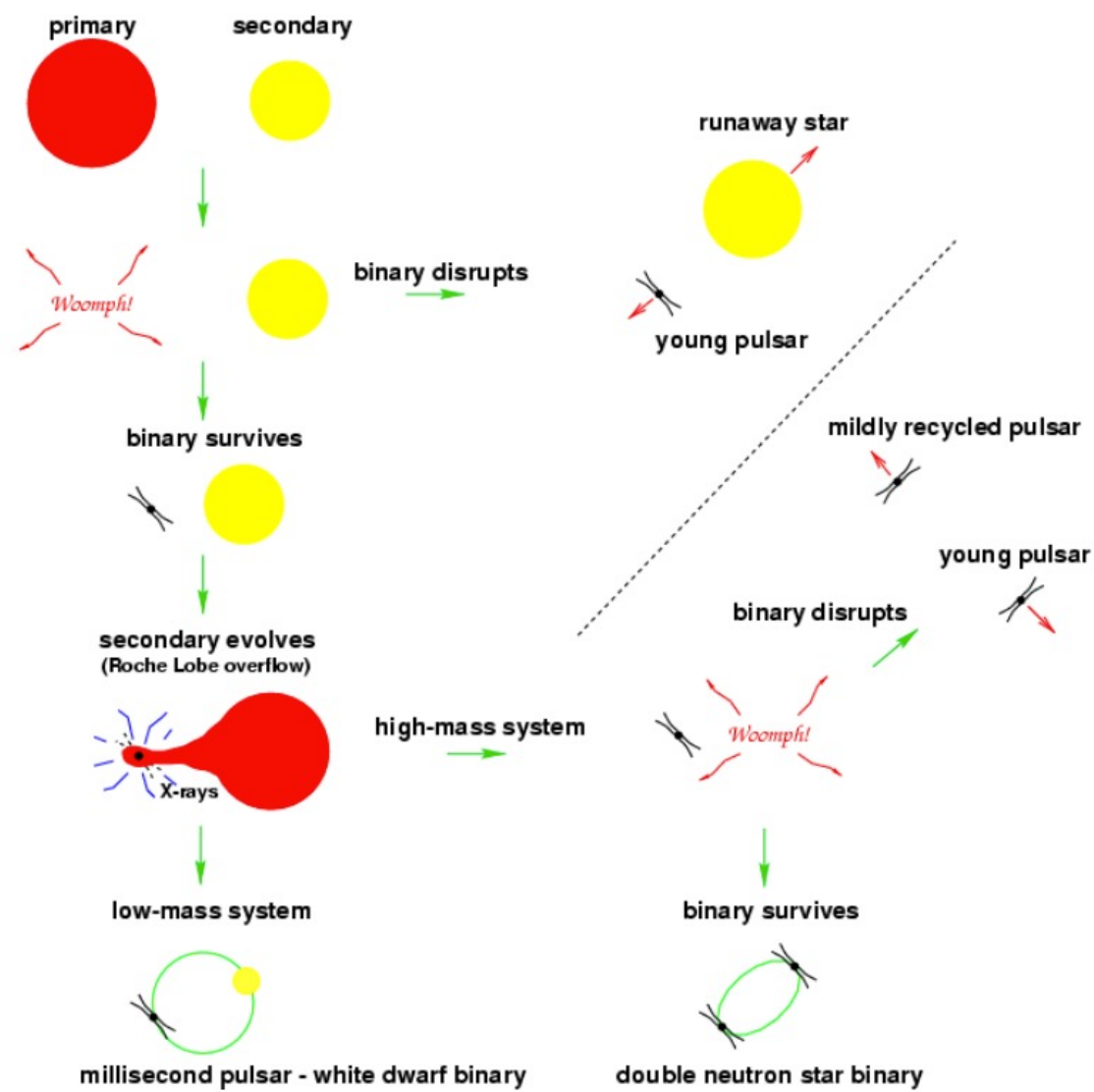
Eccentricity

- Eddy turnover time $t_{\text{eddy}} \sim \frac{r_{\text{env}}}{v_{\text{eddy}}}$
- Circularization time-scale $t_{\text{circ}} \sim t_{\text{eddy}} \frac{E_e}{E_{\text{tide}}} \sim t_{\text{eddy}} \frac{m_{\text{core}}}{m_{\text{env}}} \left(\frac{m_{\text{core}}}{M} \right)^2 \left(\frac{a}{r_{\text{env}}} \right)^8$
- Final eccentricity $e \propto \left(\frac{m_{\text{env}}}{m_{\text{core}}} \right)^{1/2} P^{1/3} v_{\text{eddy}} \propto m_{\text{env}}^{1/6} m_{\text{core}}^6 \propto P$

HMXB: $\dot{\Omega}_{\text{donor}} < \dot{\Omega}_{\text{orbit}} \rightarrow$ Darwin instability

LMXB: $\dot{\Omega}_{\text{donor}} > \dot{\Omega}_{\text{orbit}} \rightarrow$ Tides synchronize, then circularize

Lorimer (2008)



The puzzle

- Circularization theory predicts $e \sim 10^{-43}$
- Observed eccentricities of conventional MSPs are small, but non-zero
- Recent years eMSPs with $e \sim 0.1$ have been discovered
- These eMSPs deviate strongly from the general $e - P$ relation within a narrow range of orbital periods $P \approx 20 - 30$ days

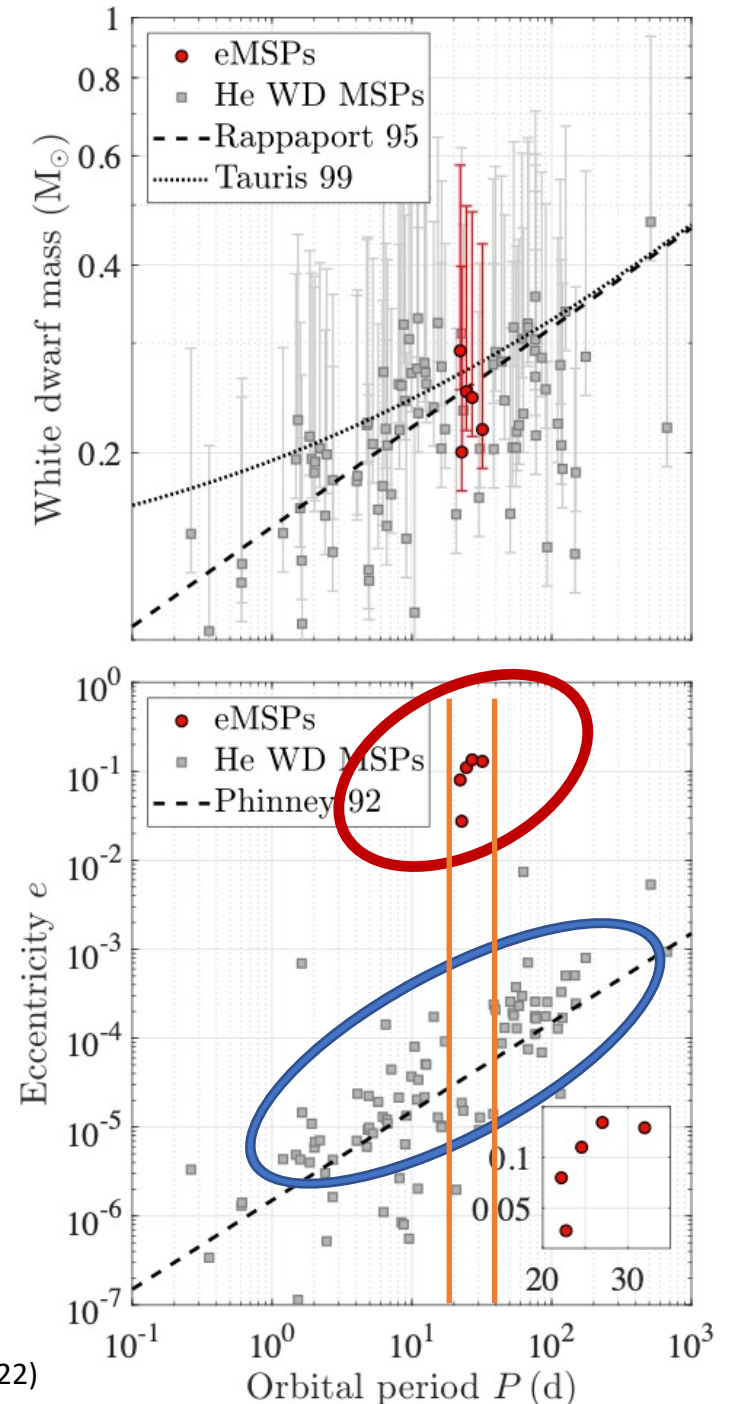
BUT

- Both MSPs and eMSPs follow the same mass-period relation

$$P \propto \left(m_{\text{core}}^{-1/3} r_{\text{env}} \right)^{3/2} \propto m_{\text{core}}^{25/4} \approx m_{\text{core}}^6$$

suggesting their common origin (RLOF of an RGB star with shell burning)

- Although residual eccentricities of MSPs can be explained (fluctuation-dissipation theorem; Phinney 1992), the existence of eMSPs poses a significant challenge



Alternative formation channels for eMSPs

Rotationally delayed accretion-induced collapse

Freire & Tauris (2014)

Predictions

- Low masses for the pulsar (1.22 – 1.31 Msun)
- Small spatial velocities for the binary

EXCLUDED (?)

Neutron star to strange quark star

Jiang et al. (2015)

Predictions

- Requires high densities \rightarrow NS > 1.8 Msun

EXCLUDED (?)

Interactions with circumbinary disk

Antoniadis (2014)

Predictions

- Pulsar masses and systemic velocities similar to conventional MSPs
- Maximum P_{orb} for eMSPs corresponds to maximum mass for He WDs that undergo flashes
- Maximum value for the eccentricity
- Smaller eccentricities possible within P_{orb} gap

Observations

- Mass measurements (e.g. PSR J1946+3417, PSR J2234+0511) seem to contradict RD-AIC & strange star scenarios
- CBD scenario so far seems to be consistent with all observations

Eccentricity maintained by convection

Fluctuation-dissipation theorem: turbulent density fluctuations in the donor's convective envelope prevent perfect circularization of the orbit (Phinney 1992)

$$e \propto \left(\frac{m_{\text{env}}}{m_{\text{core}}} \right)^{1/2} P^{1/3} v_{\text{eddy}} \propto m_{\text{env}}^{1/6} m_{\text{core}}^6 \propto P.$$

This explains well the observations for conventional MSPs except for the anomalous eMSPs, which cluster at $P \approx 20 - 30$ days

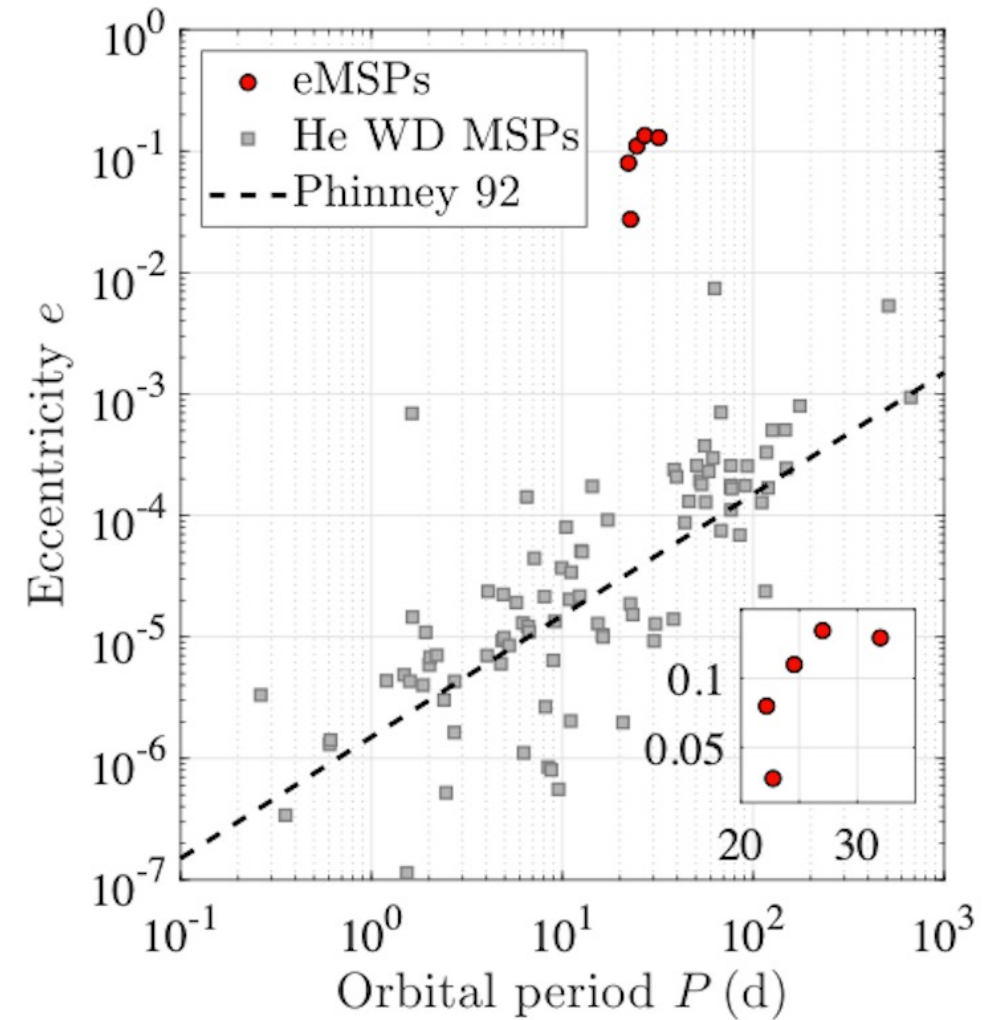
Resonance: Orbital period equals the eddy's turnover time ($P \sim t_{\text{eddy}}$). For RGB stars $t_{\text{eddy}} \approx 25$ days, exactly where the eMSPs are being found

Ansatz

At resonance eddies do not randomly change direction. Instead they form long-lived vortices generating a quadrupole moment that oscillates coherently and not stochastically \longrightarrow supported by 3D simulations of

rotating RGB stars when $Ro \equiv \frac{P_{\text{spin}}}{t_{\text{eddy}}} \sim 1$

Ginzburg & Chiang (2022)



Eccentricity enhancement by resonant convection

Results (in a nutshell)

- ✓ Assuming the eddies coherently perturb the orbit over t_{circ} , the eccentricity at resonance is enhanced by a factor of

$$\frac{e^{res}}{e} = \frac{v_e^{res}}{v_e} \sim \left(\frac{t_{circ}}{t_{eddy}}\right)^{1/2} \sim \left(\frac{t_{nuc}}{P}\right)^{1/2} \approx 3 \times 10^3$$

On Hayashi track, H^- opacity depends strongly on metallicity; Different compositions lead to $\frac{\Delta T_{eff}}{T_{eff}} \approx 0.3$ (Kippenhahn et al. 2012)

Since $t_{eddy} \sim \left(\frac{m_{env}}{\sigma T_{eff}}\right)^{1/3} \rightarrow \frac{\Delta P}{P} \approx 0.4$

- ✓ The spread in observed eMSPs orbital periods can be explained from variations in t_{eddy} due to different metallicities
- ✓ These variations broaden the range of resonant P but also leave systems out of resonance

Ginzburg & Chiang (2022)

