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# On the nature of the X-ray corona of black hole binaries in the hard state

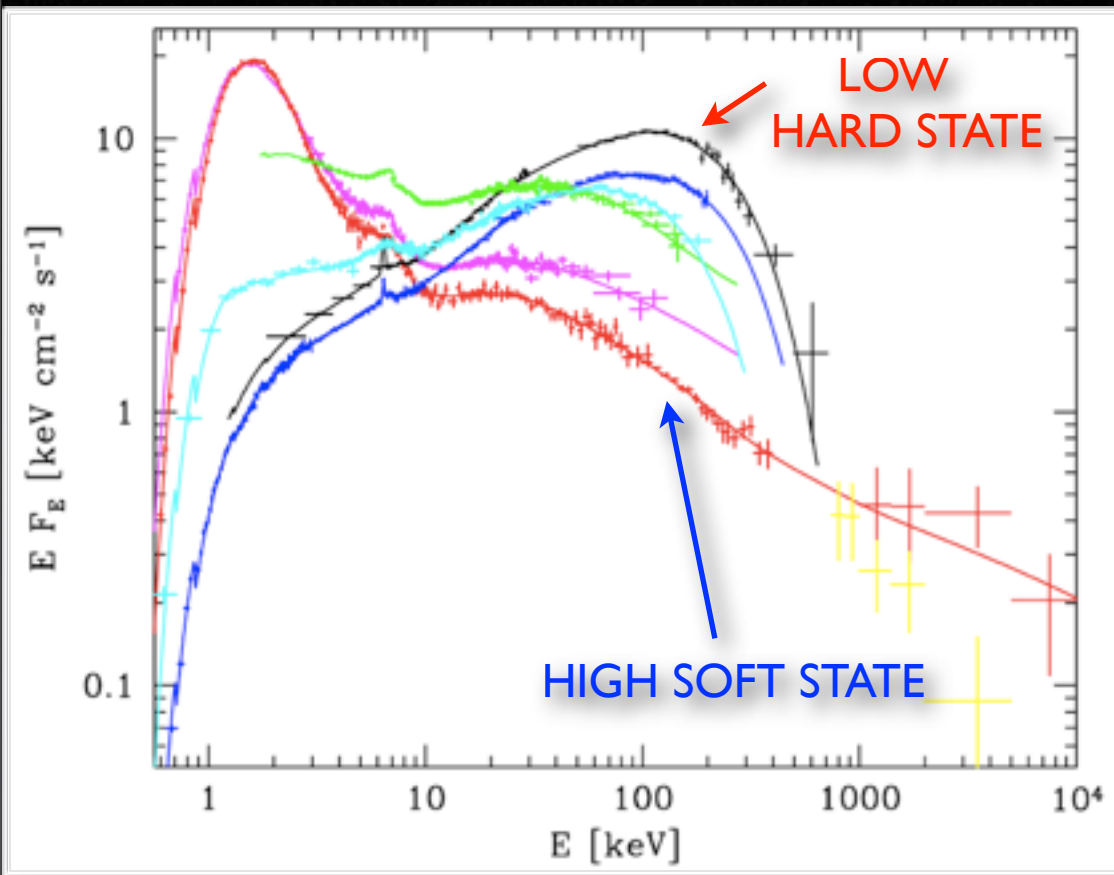
Julien Malzac (CESR/CNRS, Toulouse, France)

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## Outline

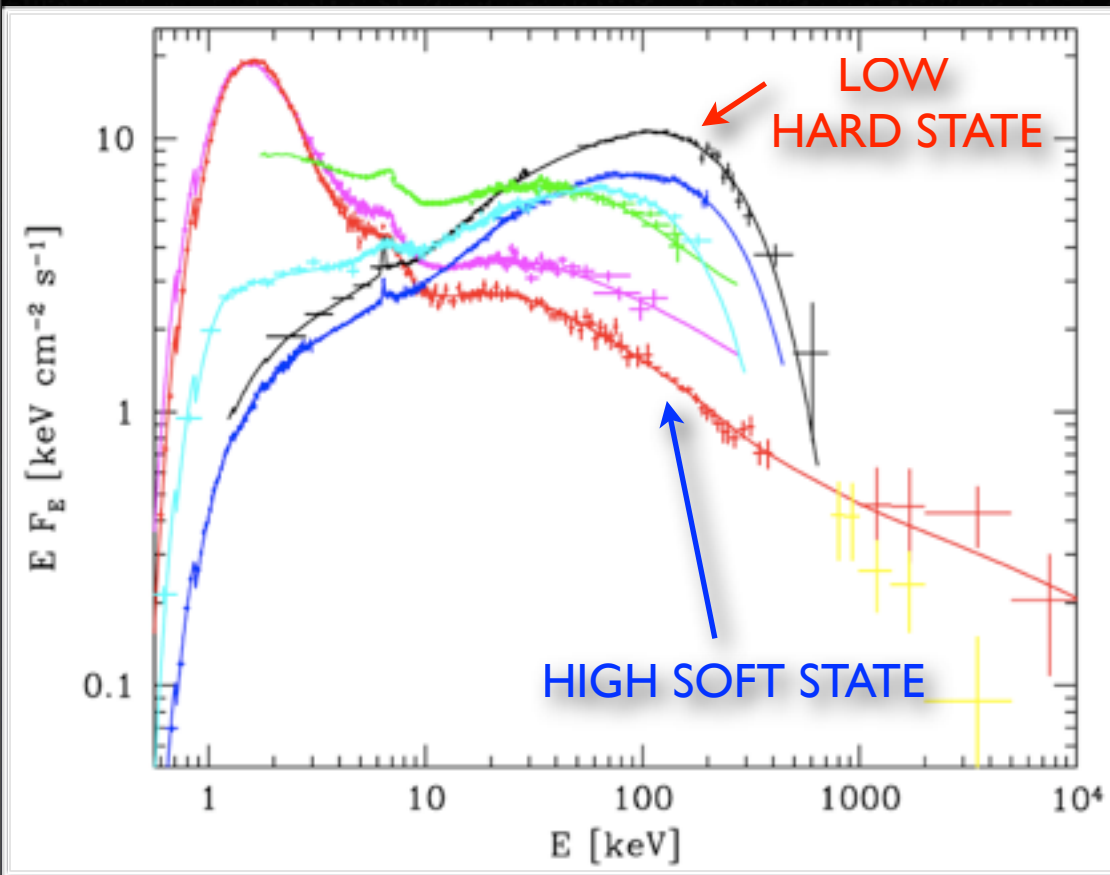
1. Black hole X-ray spectral states: observations and models
2. A numerical kinetic/radiation model for state transitions
3. Comparisons to spectra of Cygnus X-1 and GX 339-4
4. Applicability of hot accretion flow models to bright hard state BHBs
5. Toward a multi-zone corona model

# High energy emission of Cygnus X-1

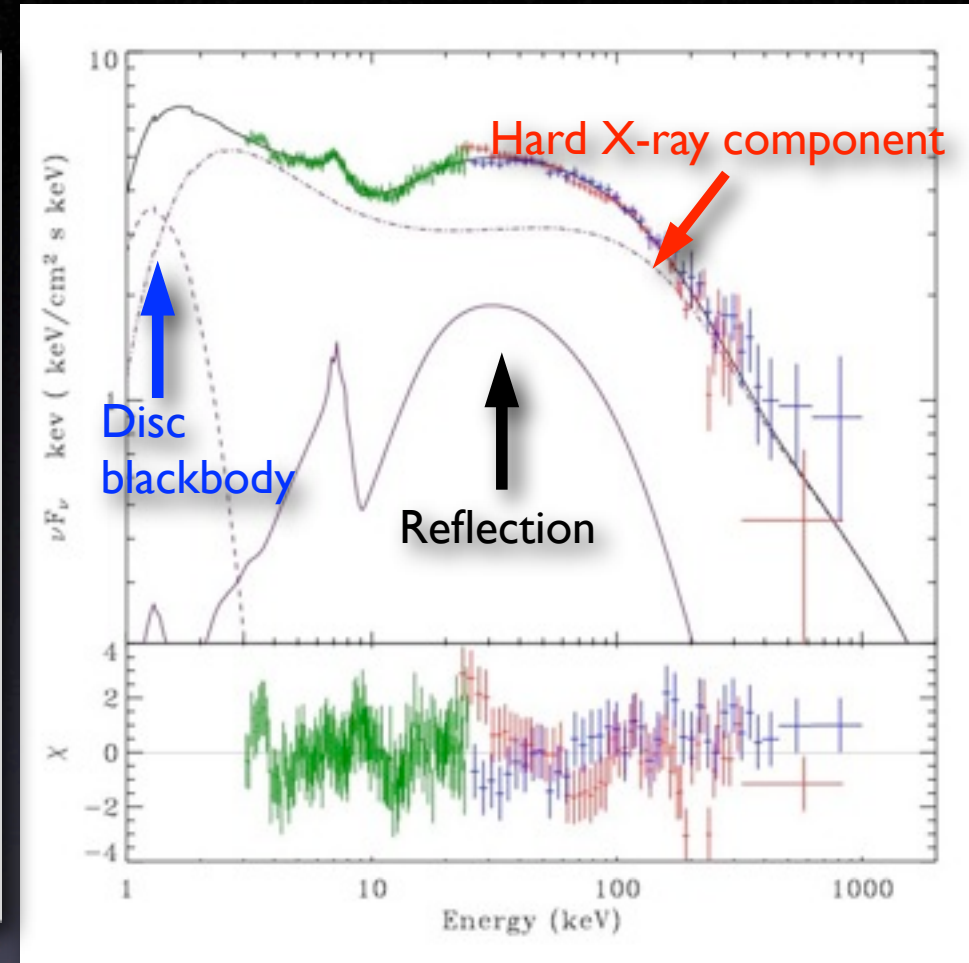


Zdziarski et al 2003

# High energy emission of Cygnus X-1



Zdziarski et al 2003



Malzac et al. 2006

**LOW HARD STATE:** (compact radio jet)

disc blackbody and reflection: weak / Corona: THERMAL Comptonisation

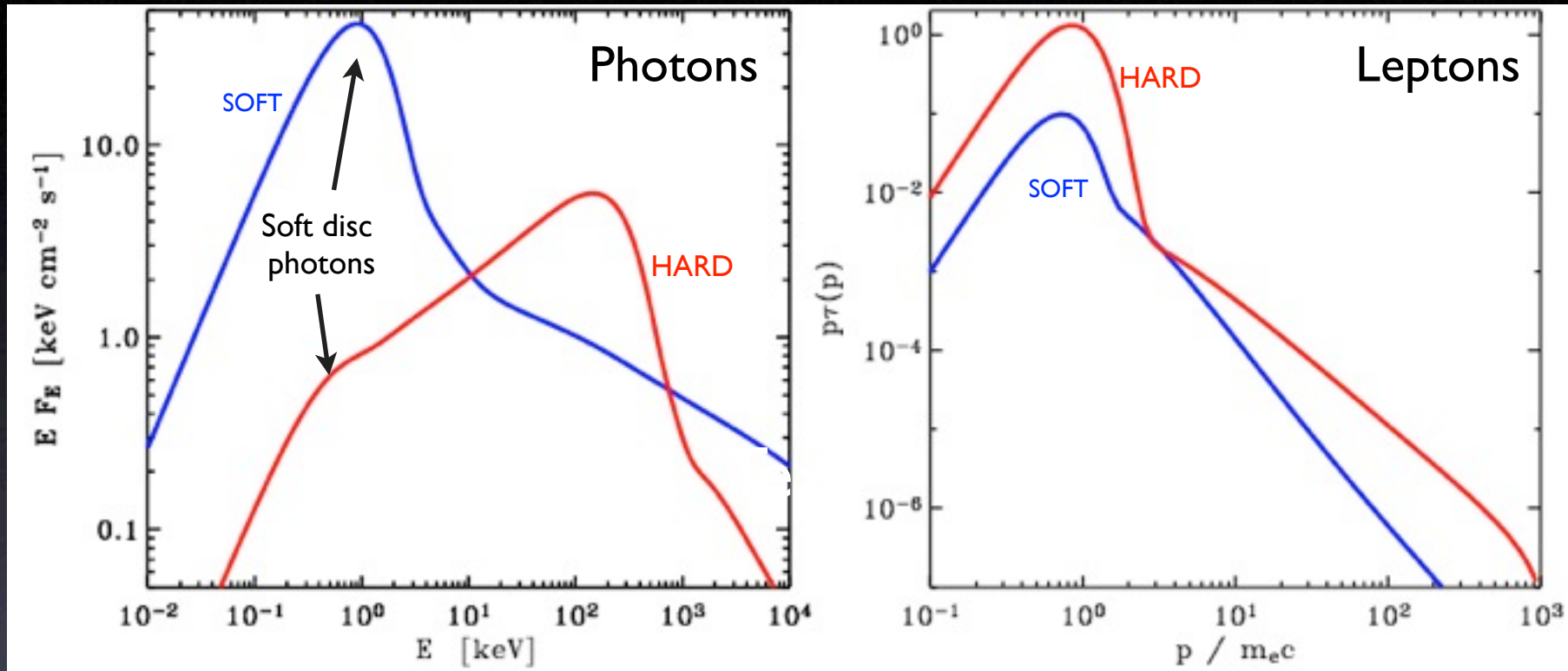
**HIGH SOFT STATE:**

disc blackbody and reflection: strong / Corona: NON-THERMAL Comptonisation



# Hybrid thermal/non-thermal comptonisation models

(Coppi 1999; Gierlinski et al. 1999, Zdziarski ..., Done ...)



- Comptonising electrons have similar energy distribution in both states:  
Maxwellian+ non-thermal tail

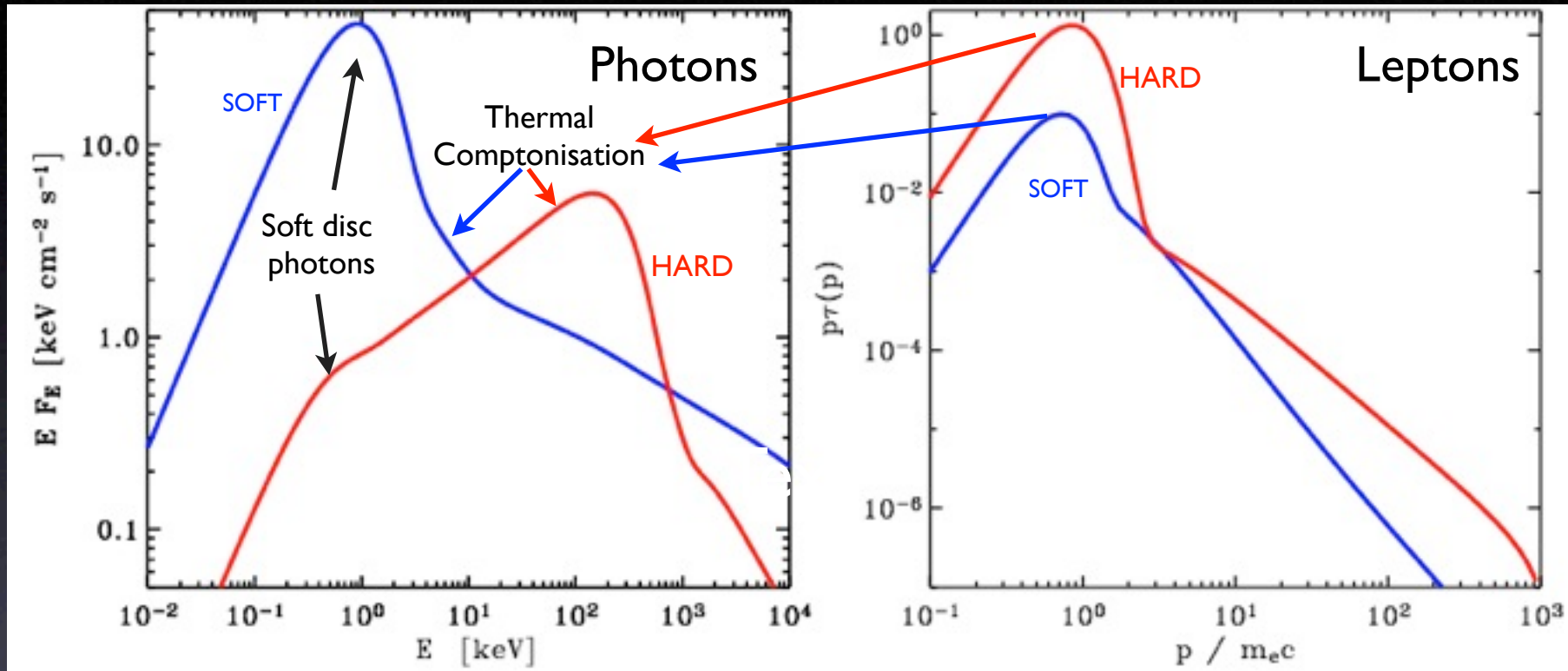
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**SOFT STATE:**  $kT \sim 10-50 \text{ keV}$ ,  $\tau_T \sim 0.1-0.3$ : Inverse Compton by non-thermal electrons dominates

- Lower temperature of corona in soft state possibly due to radiative cooling by soft disc photons

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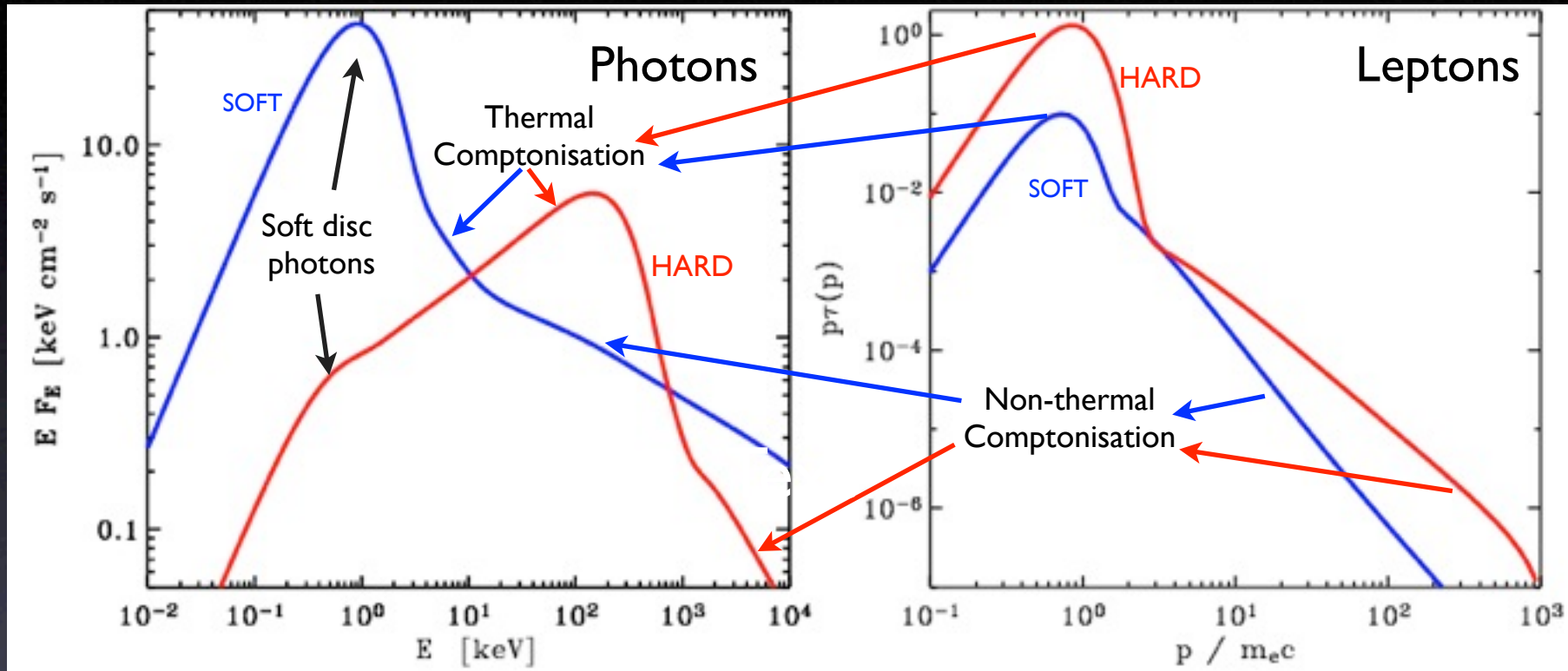
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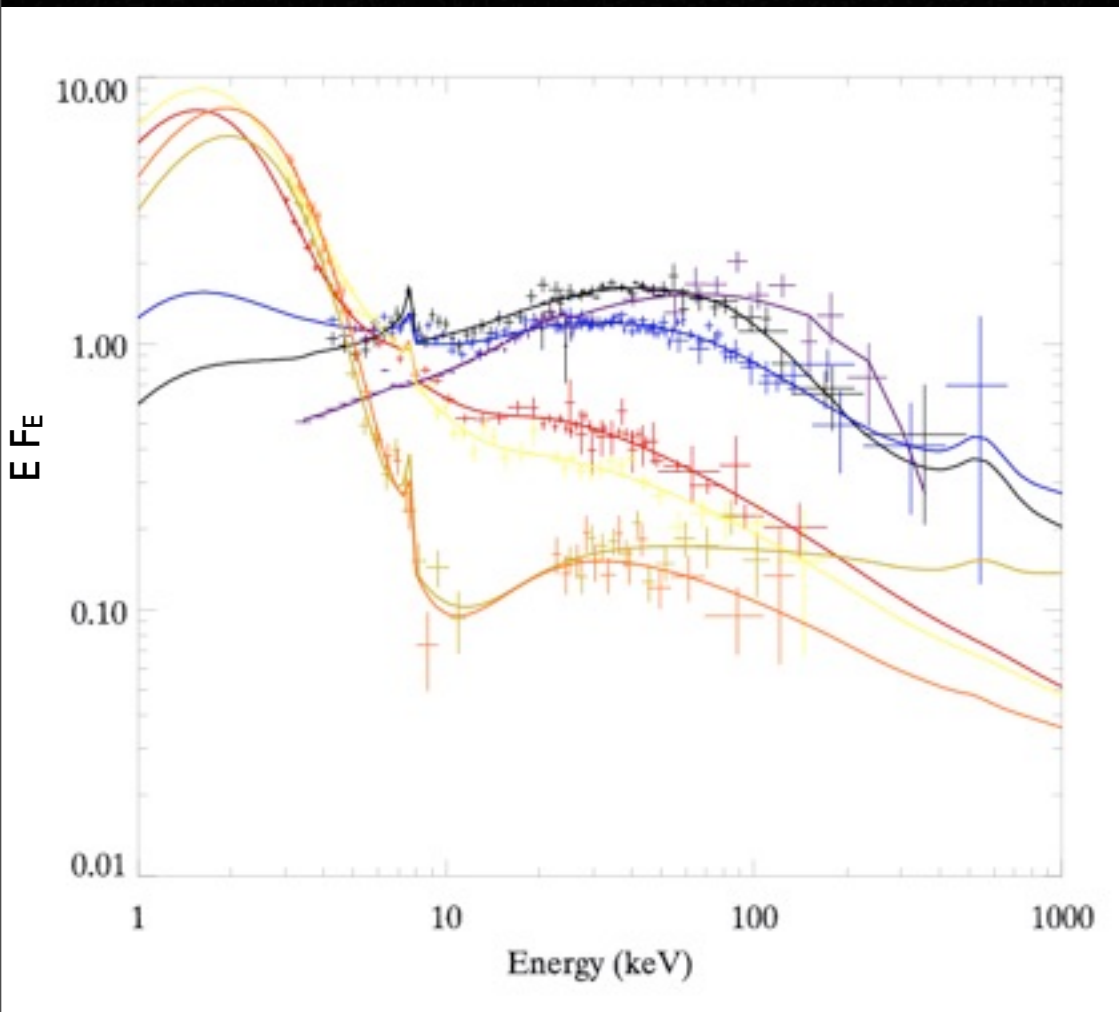
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# GX 339-4 during the 2004 state transition



INTEGRAL

Del Santo, Malzac, Jourdain, Belloni, Ubertini, MNRAS, 2008

see also Joinet et al. (2007), Belloni et al. (2006),

- Smooth transition from thermal to non-thermal Comptonisation
- Fits with hybrid thermal/non-thermal models (EQPAIR) during the Hard to Soft transition:
  - softening driven by dramatic cooling of the coronal electrons by soft disc photons

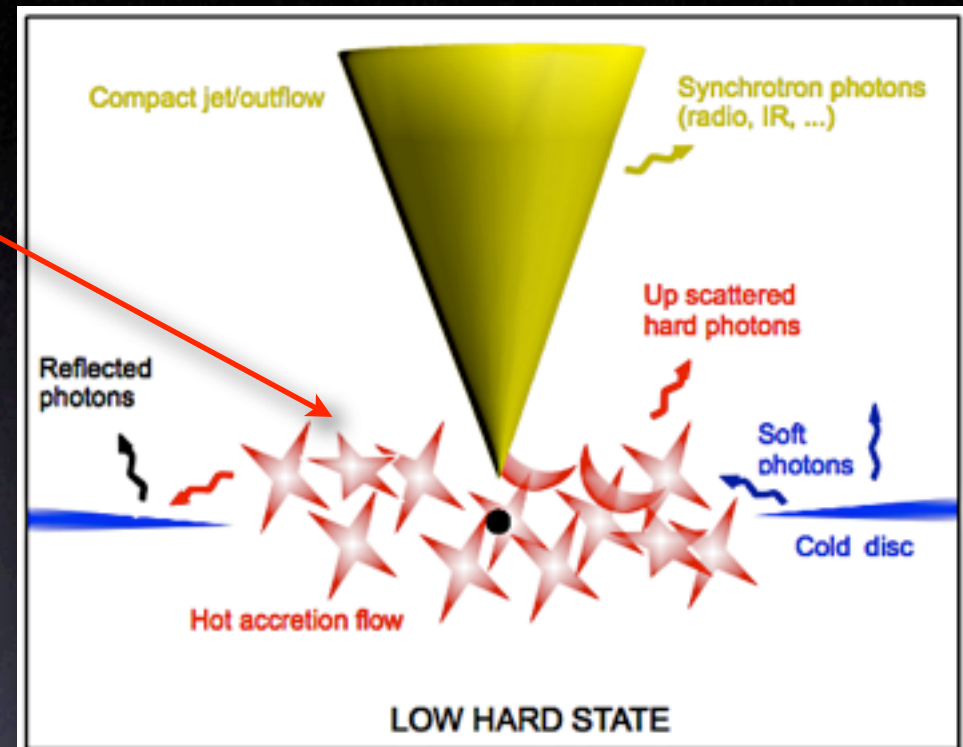
# Standard picture: truncated disc model

## LOW HARD STATE

cold disc truncated at  $\sim 100-1000 R_g$   
+ hot inner accretion flow

$\Rightarrow$  Thermal Comptonisation  
in the hot ( $10^9$  K) plasma

(Shapiro, Lighman & Eardley 1976; Rees et al. 1982;  
Narayan & Yi 1994, Abramowicz et al. 1995, Esin et al.  
1997, Yuan & Zdziarski 2004, Petrucci et al. 2010...)



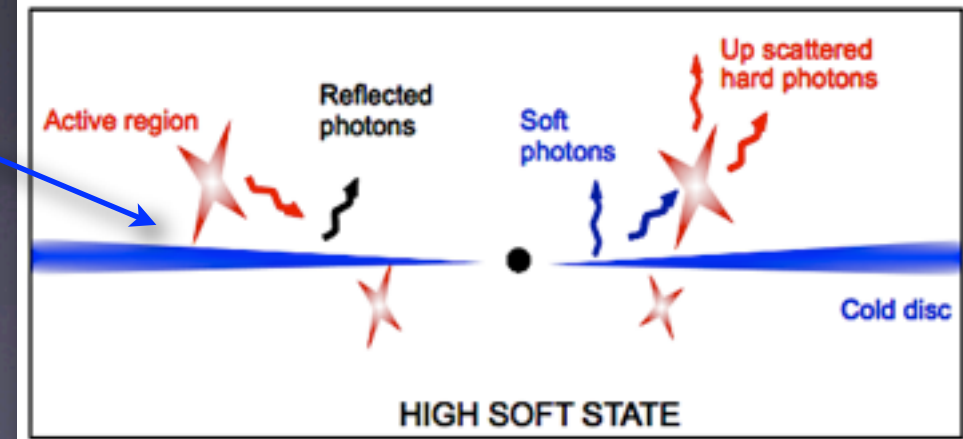
## HIGH SOFT STATE

cold geometrically thin disc  
down to the last stable orbit  
+ weak non-thermal corona

$\Rightarrow$  dominant thermal disc emission

+ non-thermal Comptonisation

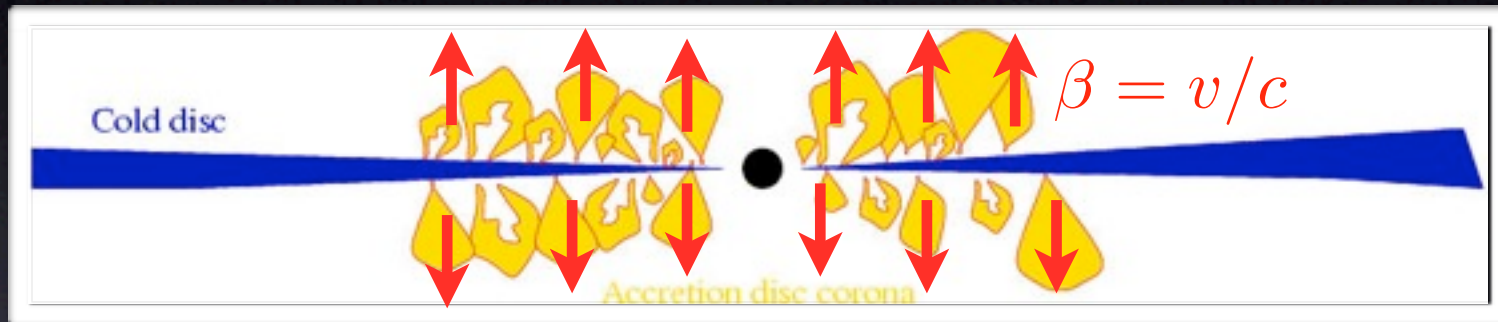
(Shakura & Sunyaev 1973, Galeev et al. 1979, Coppi 1999)





# Alternative models for the hard state

- Accretion disc corona outflowing with mildly relativistic velocity above a cold (i.e. non-radiating) thin disc



(Beloborodov 1999; Malzac Beloborodov & Poutanen 2001)

## ● X-ray Jet Models

(Markoff et al. 2001,2005; Reig et al. 2003; Giannios et al. 2004; Kylafis et al. 2008)

# Global energy budget in Cyg X-1

● Jet powered nebula:  $P_j \simeq L_X \simeq 2 \times 10^{37} \text{ erg s}^{-1}$  (Gallo et al. 2005, Russell et al. 2007)





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➔ accretion proceeds efficiently in the hard state

➔ cannot be strongly advection dominated

➔ not enough power to eject corona with  $\tau_T > 1$   
to infinity with relativistic speed

➔ X-ray corona  $\neq$  Jet



# BELM: a code to model radiation and kinetic processes in the corona

➔ Evolution of electrons and photon energy distributions in a fully ionised, magnetised plasma (radiation, acceleration and Coulomb processes)

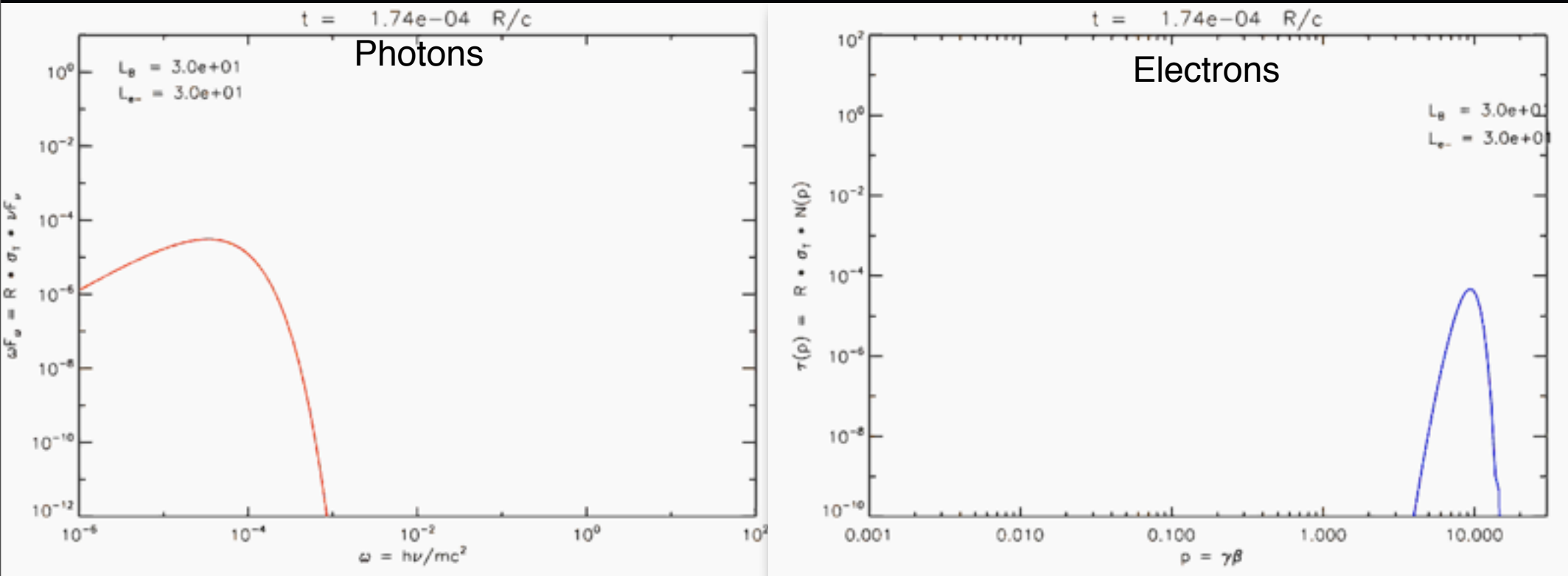
● Solve coupled time-dependent kinetic equations (one zone) for leptons and photons (no assumption on the shape of the electron distributions)

● Compton, Synchrotron emission and absorption, e-e and e-p Coulomb,  $e^+e^-$  pair production/annihilation, e-p bremsstrahlung

(Belmont, Malzac & Marcowith, A&A 2008)

# The Synchrotron boiler

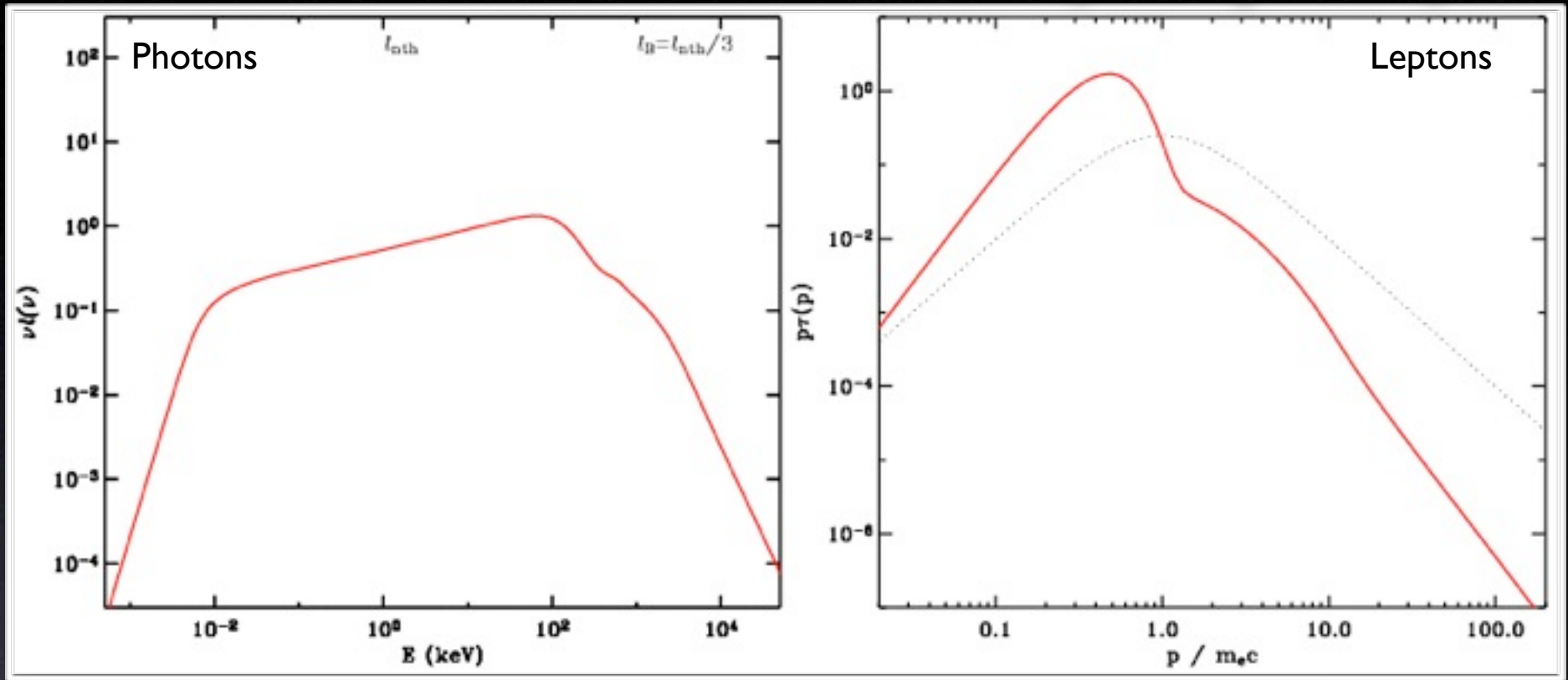
(Ghisellini, Guilbert and Svensson 1988)



- Electrons injected with  $\gamma = 10$  in an empty (but magnetised) region  
Synchrotron self-Compton emission
- High energy  $e^- \rightarrow$  synchrotron photons  $\rightarrow$  self-absorbed by lower energy  $e^-$   
 $\rightarrow$  transfer of energy between particles
  - $\rightarrow$  'thermalizing' effect on the electron distribution
  - $\rightarrow$  At steady state: hybrid thermal/non thermal lepton distribution

(Belmont, Malzac & Marcowith, A&A, 2008)

# Pure non-thermal SSC models (steady state)

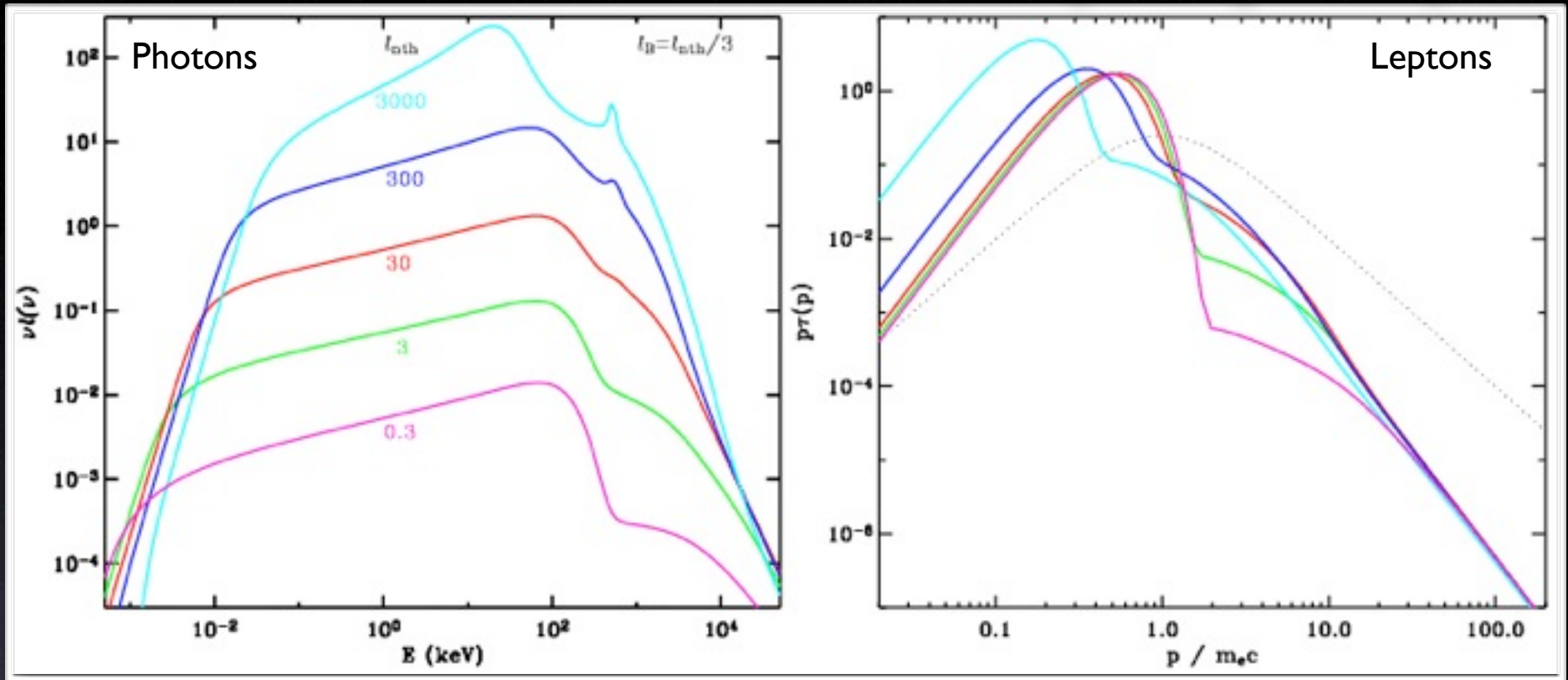


- ➊ Magnetic field  $B$  at  $\sim$ equipartition with radiation,  $l_B = (\sigma_T/m_e c^2) R B^2 / (8\pi)$
- ➋ Continuous POWER-LAW electron injection  $\Gamma_{inj}=3$ ,  $l_{nth} = (\sigma_T/m_e c^3) L/R$
- ➔ Cooling and thermalisation through synchrotron self-Compton + e-e Coulomb
- ➔ Equilibrium distribution: Maxwellian+ non-thermal tail
- ➔ **spectra look like hard state !**

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# Effect of external soft photons

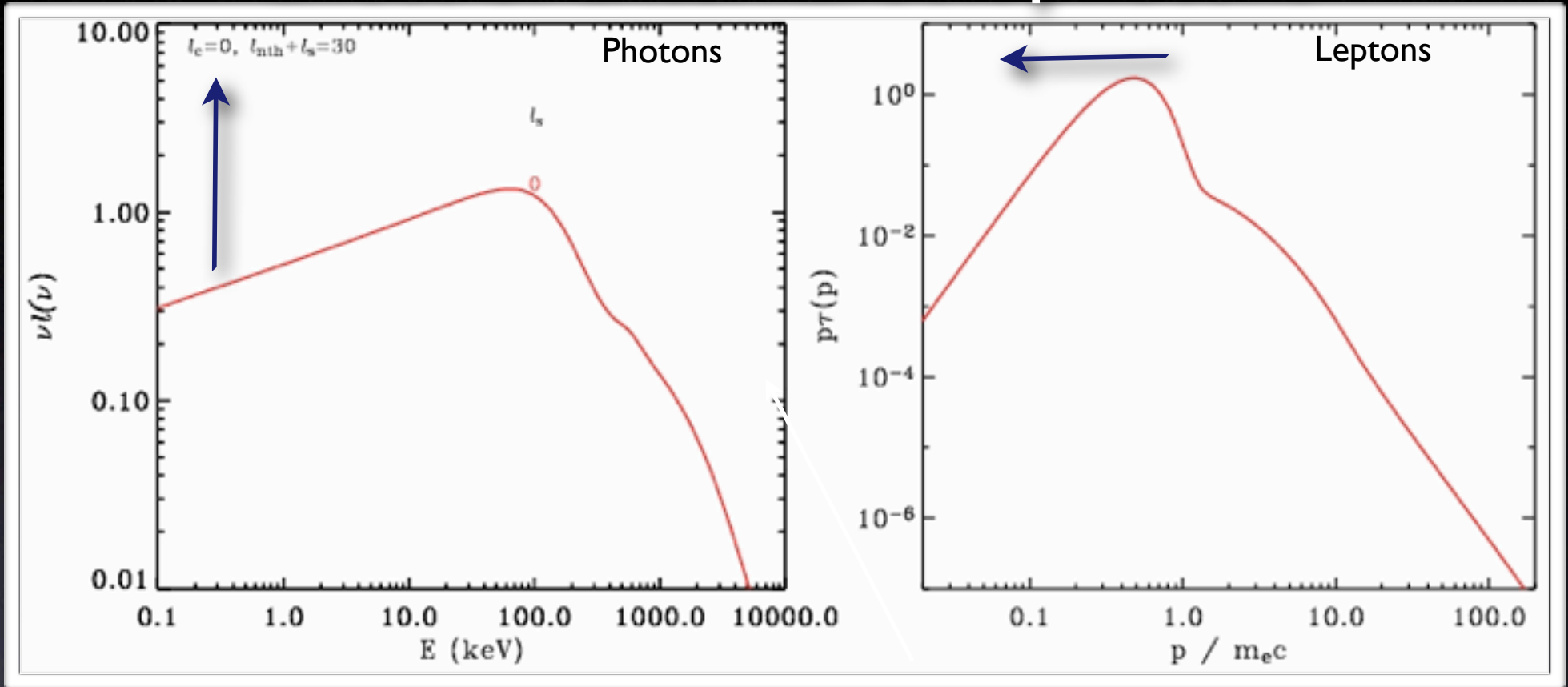


- Add soft thermal photons:
  - ➔ temperature of Maxwellian electrons decreases
  - ➔ Compton emission increasingly dominated by non-thermal electrons
  - ➔ looks like a state transition!

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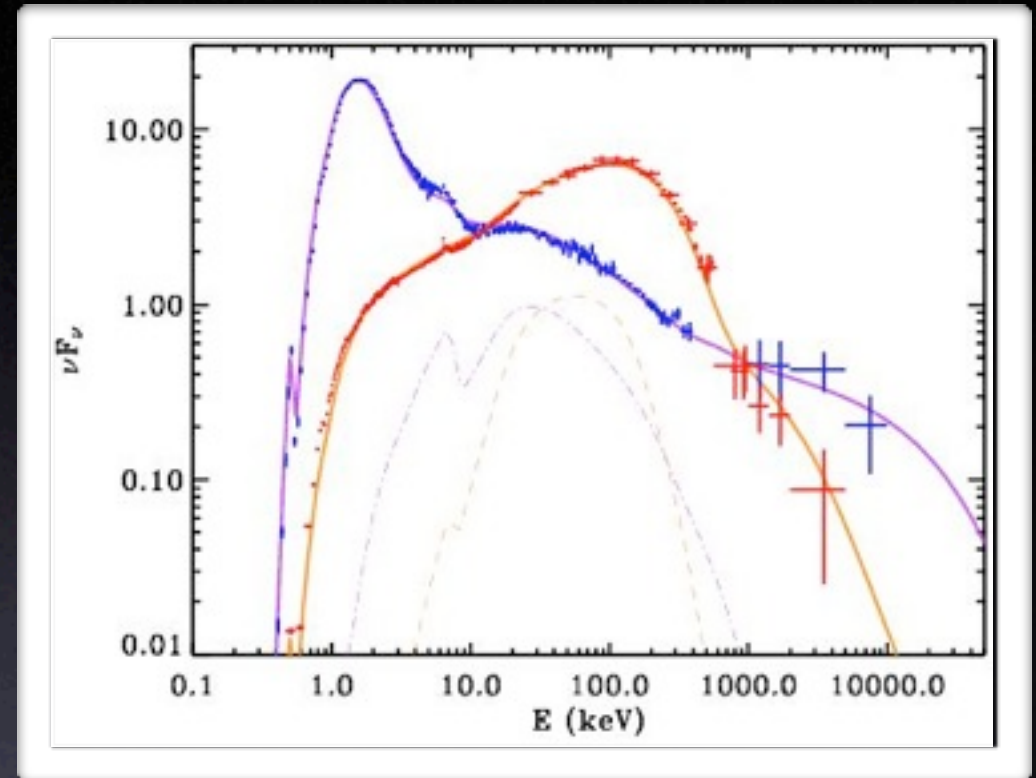


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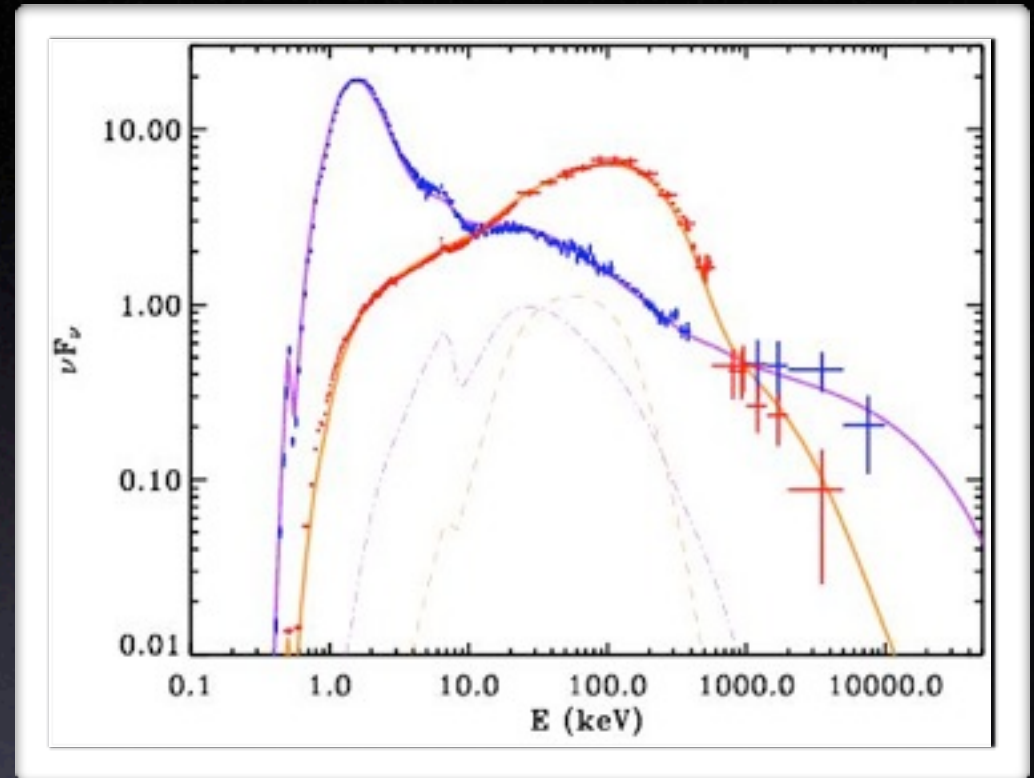


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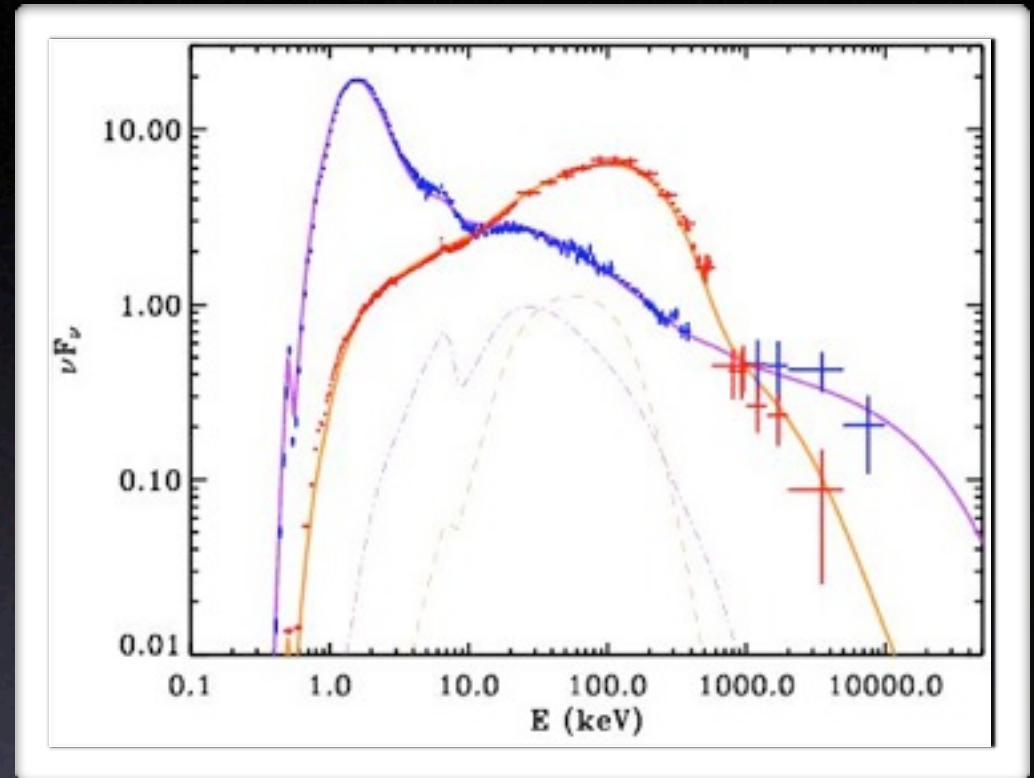
- Both states consistent with pure non-thermal acceleration models
- Different coronal temperatures due to more cooling by thermal disc photons in Soft state



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If B is large:

- ➔ non-thermal electrons generate too much synchrotron
- ➔ Maxwellian electrons are too cold
- ➔ **weak (i.e strongly sub-equipartition) magnetic field**
- ➔ **corona unlikely to be powered by magnetic field**

(Malzac & Belmont 2009 ; Poutanen & Vurm 2009)



# Model with hot protons

In addition to non-thermal acceleration we now assume that electrons are heated through Coulomb interactions with a population of hot thermal protons (two-temperature plasma):

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- Temperature of hot protons in hard state:

$$T_i < 2 \cdot 10^{10} \text{ K or } T_i/T_e < 10$$

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 $T_i < 2 \cdot 10^{10} \text{ K}$  or  $T_i/T_e < 10$   
➔ **proton temperature much lower than standard two-temperature accretion disc solutions**
- Similar constraints on  $B$  and  $T_i$  obtained for GX339-4 in a bright hard state (Droulans et al. 2010)



# Can hot accretion flows explain the bright hard state sources?

● In the context of alpha discs, (i.e.  $Q_{\text{vis}} = -\alpha P_{\text{gas}} R \frac{d\Omega}{dr}$ ),

there is no hot flow solutions with  $\tau_{\text{T}} \geq 1$ : cooling is too strong.

➔ standard ADAF solution cannot be applied

● A possible fix:

1) Assume  $P_{\text{mag}} \geq P_{\text{gas}}$

2) Modified viscosity law:  $Q_{\text{vis}} = -\alpha(P_{\text{gas}} + P_{\text{mag}})R \frac{d\Omega}{dr}$

➔ solutions with  $\tau_{\text{T}} \geq 1$ ,  $T_{\text{i}}/T_{\text{e}} \sim 2 - 10$ ,  $P_{\text{mag}}/P_{\text{gas}} \sim 2$

(e.g. Oda et al 2010, Bu et al 2009, Fragile & Meier 2009, Malzac et al in prep)

- Hot accretion flow solutions
- Accretion disk coronae → strong magnetic field
- MHD jet models

but...

- Non-thermal high energy excess → weak magnetic field

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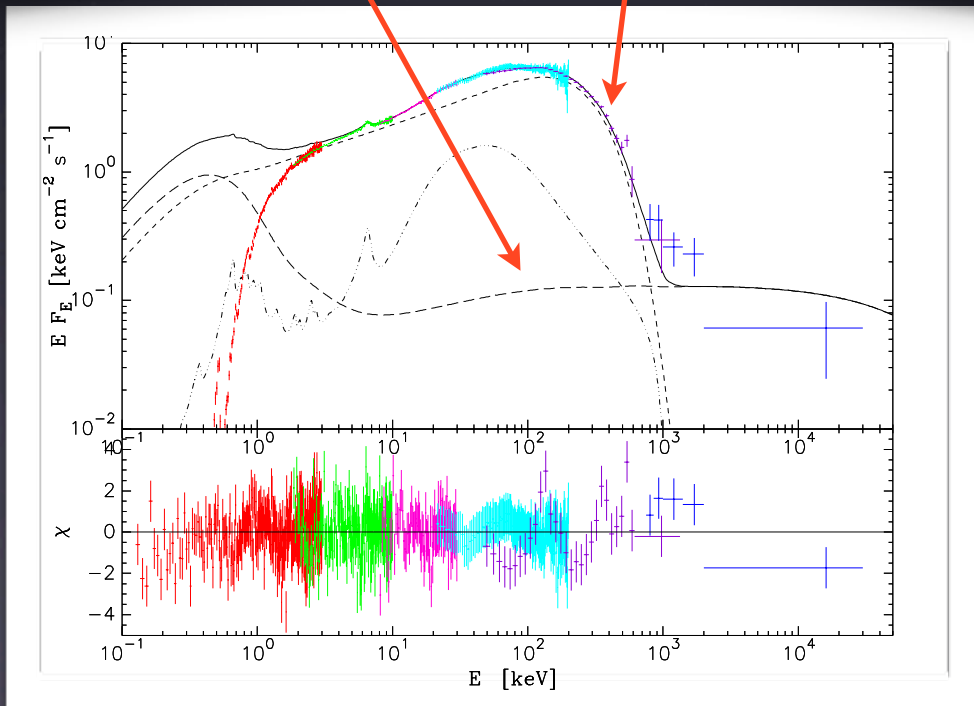
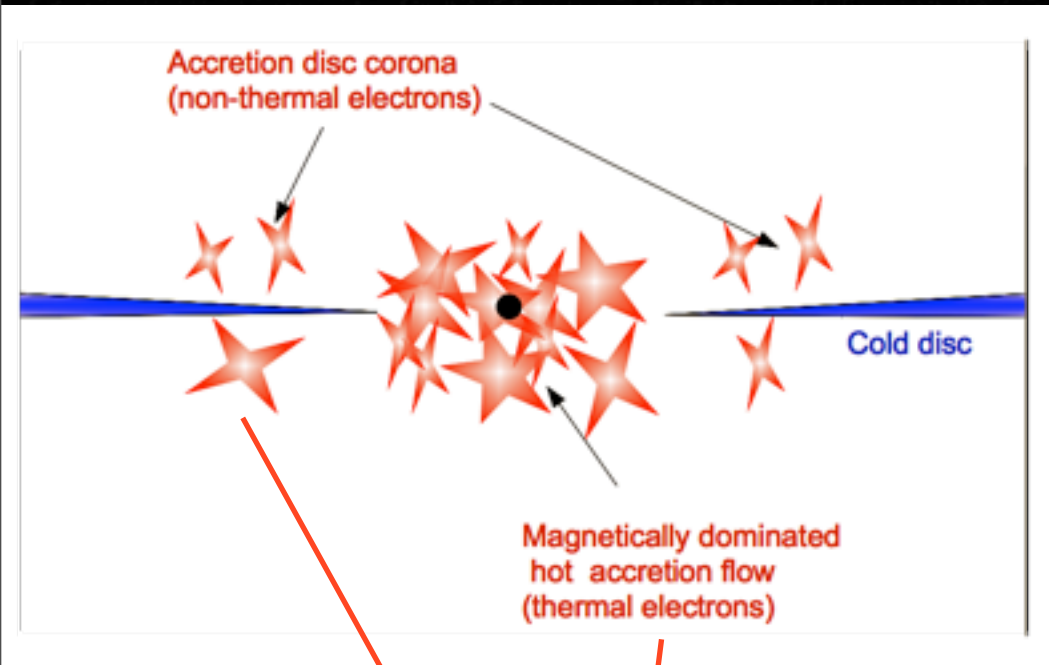
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Constraint of low B removed if thermal and non-thermal  
Comptonisation produced in different locations

→ multi-zone corona ?

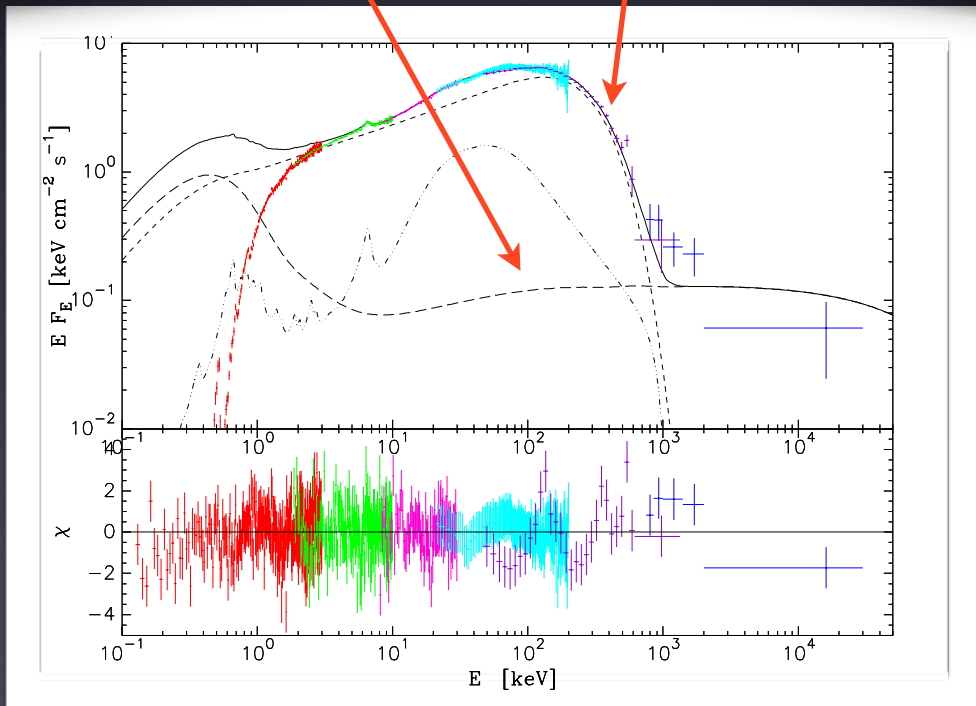
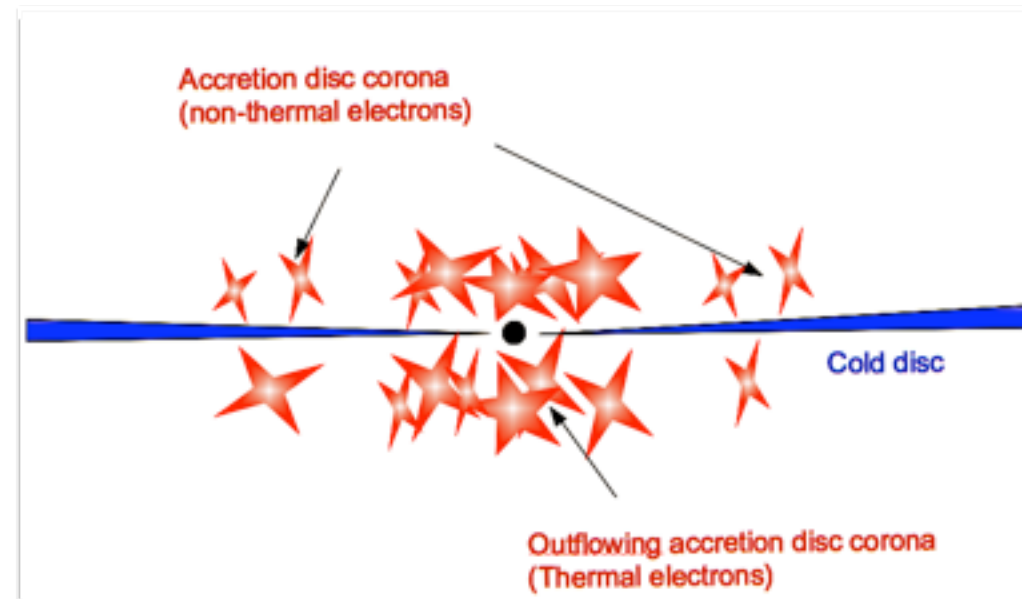
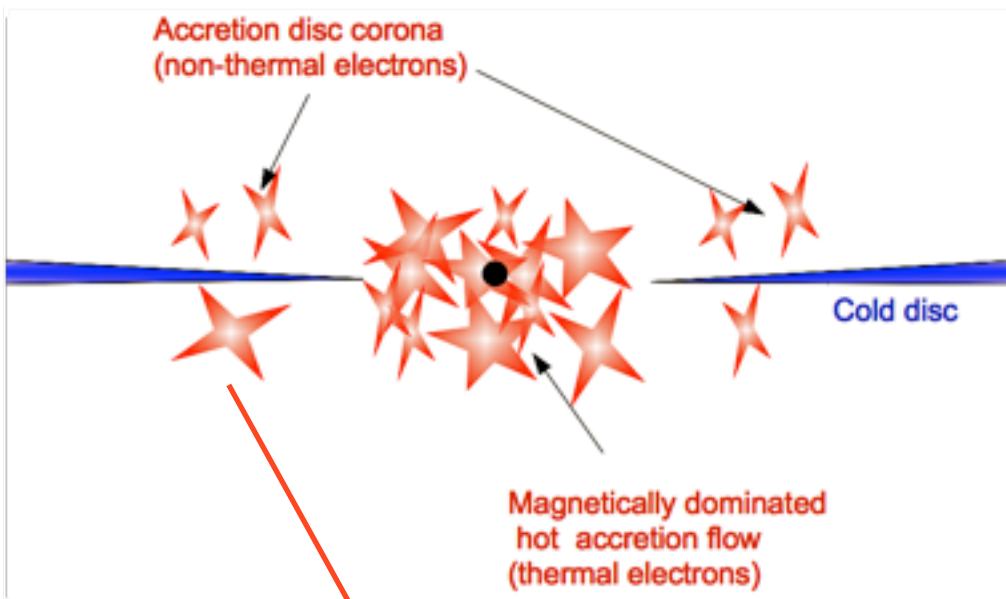
# A two-component model for the LHS



- Thermal comptonisation component dominates hard X-ray emission
- Non-thermal component reproduces soft X-ray excess and MeV emission

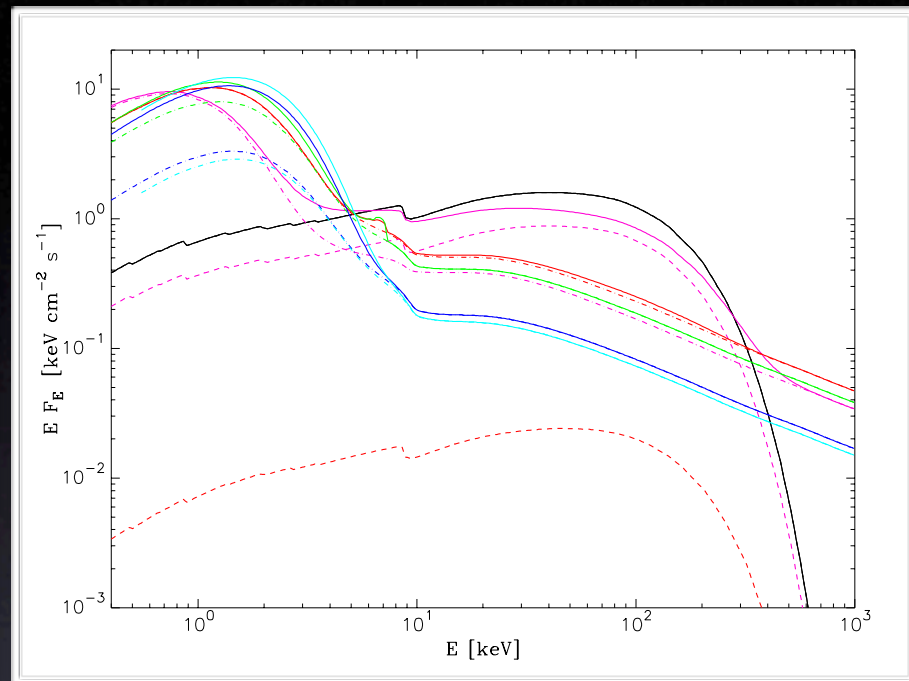
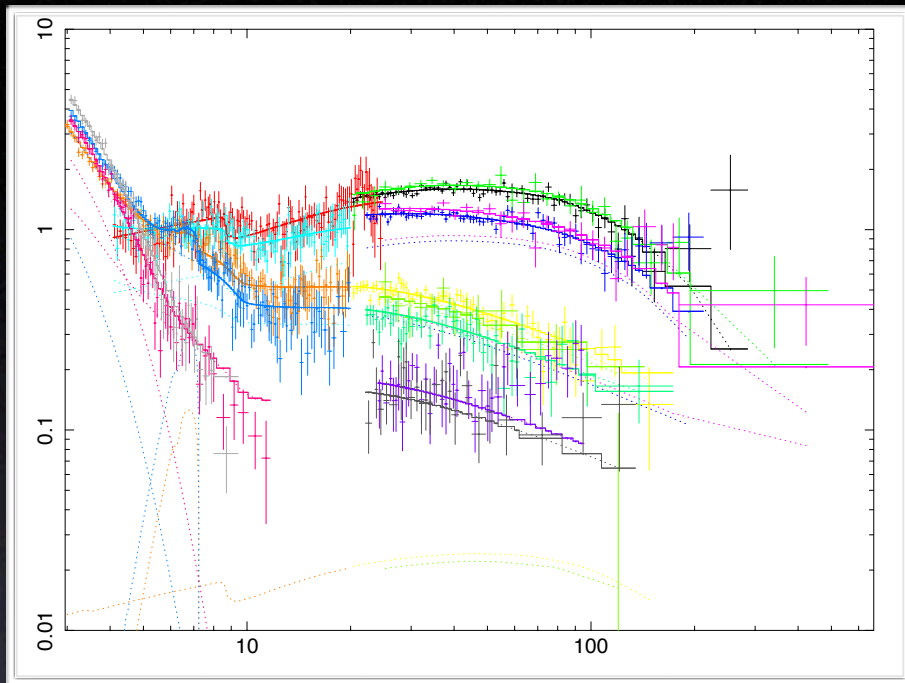


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# Spectral state transitions revisited



● INTEGRAL data from GX339-4 during state transition consistent with a corona model with 2 zones (pure thermal and non-thermal).

● Shapes of thermal and non-thermal components are constant. Only temperature of disc blackbody and normalisations vary during transition.



# Conclusions:

- We still do not know what the corona is ...
- In the best documented sources, none of the 'usual' corona models really fits the data
- However, magnetically dominated hot flow models seem promising for bright hard state sources.
- Magnetic field likely to be strong, effects on
  - accretion flow dynamics
  - particle thermalisation / cooling
  - radiation
- If so the structure of the corona appears complex: multi-zone models appear required